

UNIVERSITY OF SWAZILAND

144

FACULTY OF SCIENCE

DEPARTMENT OF PHYSICS

SUPPLEMENTARY EXAMINATION 2011/2012

TITLE OF PAPER : ELECTROMAGNETIC THEORY I

COURSE NUMBER : P331

TIME ALLOWED : THREE HOURS

**INSTRUCTIONS : ANSWER ANY FOUR OUT OF FIVE
QUESTIONS.
EACH QUESTION CARRIES 25 MARKS.**

**MARKS FOR DIFFERENT SECTIONS ARE
SHOWN IN THE RIGHT-HAND MARGIN.**

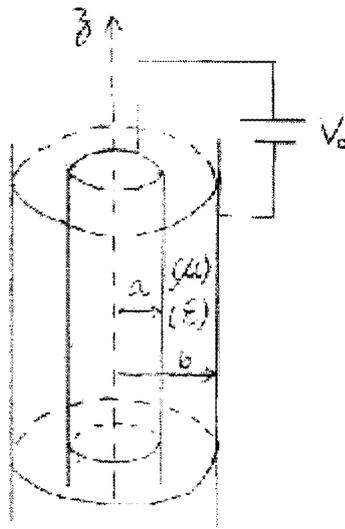
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P331 Electromagnetic Theory I

Question one

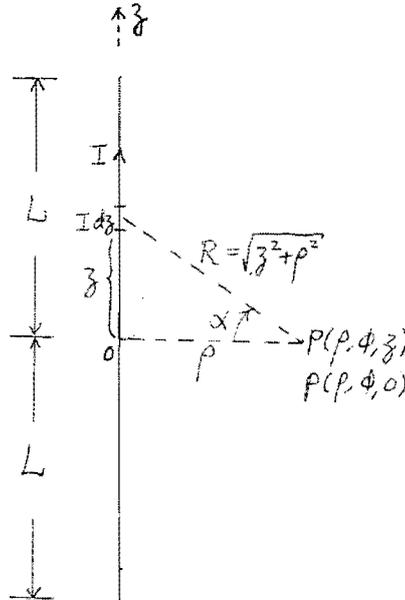
A very long straight coaxial cable has its central axis aligned with z -axis, its inner and outer conducting cable's cross-sectional radius are a & b respectively and in-between the cables is a insulating layer of permittivity ϵ , as shown below :



- (a) From $\nabla^2 f(\rho) = 0$, i.e., $\frac{d}{d\rho} \left(\rho \frac{df(\rho)}{d\rho} \right) = 0$, with boundary conditions of $f(\rho = b) = 0$ and $f(\rho = a) = V_0$, find the specific solution of $f(\rho)$ and show that $f(\rho) = -\frac{V_0}{\ln\left(\frac{b}{a}\right)} (\ln(\rho) - \ln(b))$. **(10 marks)**
- (b) Use $\vec{E} = -\vec{\nabla} f$ to find the electric field \vec{E} in the insulating layer, i.e., $a \leq \rho \leq b$. **(4 marks)**
- (c) From $\rho_s = \epsilon E_n$ find the electric surface charge density ρ_s on $\rho = a$ and $\rho = b$ surfaces respectively. Then find the total charges deposited on both conducting cable surfaces (i.e., on $\rho = a$ & $\rho = b$ surfaces) of one meter length and show that they are equal and opposite. **(9 marks)**
- (d) Find the distributive capacitance of this coaxial cable system in terms of a , b & ϵ . **(2 marks)**

Question two

- (a) For the time-independent (or static) case, setting $\vec{B} = \vec{\nabla} \times \vec{A}$ and using Coulomb's gauge, i.e., $\vec{\nabla} \cdot \vec{A} = 0$, to deduce the following Poisson's equations for \vec{A} from Maxwell's equations as $\nabla^2 \vec{A} = -\mu \vec{J}$. (5 marks)
- (b) A straight thin conducting wire of length $2L$ carries a total current I along $+z$ direction as shown in the diagram below:



- (i) find the z -component of vector potential \vec{A} , i.e., A_z , due to the above given current at a point $P : (\rho, \phi, 0)$ by evaluating the following integral

$$A_z = \int_{z=-L}^L \frac{\mu_0 I dz}{4\pi R} = 2 \int_{z=0}^L \frac{\mu_0 I dz}{4\pi \sqrt{z^2 + \rho^2}} \quad \text{and show that}$$

$$A_z = \frac{\mu_0 I}{2\pi} \left(\ln \left(\frac{\sqrt{\rho^2 + L^2} + L}{\rho} \right) \right) \quad \text{(10 marks)}$$

(Hint: $d \tan(\alpha) = \sec^2(\alpha) d\alpha$ & $1 + \tan^2(\alpha) = \sec^2(\alpha)$,
 $\int \sec(\alpha) d\alpha = \ln(\sec(\alpha) + \tan(\alpha))$)

- (ii) From $\vec{B} = \vec{\nabla} \times \vec{A}$ and knowing $A_\rho = 0 = A_\phi$, find the magnetic field \vec{B} at the point P due to the above given current. In the case of $L \gg \rho$, show that magnetic field \vec{B} obtained here can be simplified to $\vec{B} = \vec{a}_\rho \frac{\mu_0 I}{2\pi \rho}$. (10 marks)

Question three

Starting with the following Maxwell's equations for a material region with parameters of $(\mu, \varepsilon, \sigma)$ as :

$$\left\{ \begin{array}{l} \vec{\nabla} \cdot \vec{E}(space, t) = 0 \quad \dots\dots (1) \\ \vec{\nabla} \cdot \vec{H}(space, t) = 0 \quad \dots\dots (2) \\ \vec{\nabla} \times \vec{E}(space, t) = -\mu \frac{\partial \vec{H}(space, t)}{\partial t} \quad \dots\dots (3) \\ \vec{\nabla} \times \vec{H}(space, t) = \sigma \vec{E}(space, t) + \varepsilon \frac{\partial \vec{E}(space, t)}{\partial t} \quad \dots\dots (4) \end{array} \right.$$

(a) setting $\vec{E}(space, t) = \vec{\tilde{E}}(space) e^{i\omega t}$ & $\vec{H}(space, t) = \vec{\tilde{H}}(space) e^{i\omega t}$, deduce the following time-harmonic Maxwell's equations :

$$\left\{ \begin{array}{l} \vec{\nabla} \cdot \vec{\tilde{E}}(space) = 0 \quad \dots\dots (5) \\ \vec{\nabla} \cdot \vec{\tilde{H}}(space) = 0 \quad \dots\dots (6) \\ \vec{\nabla} \times \vec{\tilde{E}}(space) = -i\omega\mu \vec{\tilde{H}}(space) \quad \dots\dots (7) \\ \vec{\nabla} \times \vec{\tilde{H}}(space) = (\sigma + i\omega\varepsilon) \vec{\tilde{E}}(space) \quad \dots\dots (8) \end{array} \right. \quad \text{(3 marks)}$$

(b) consider the space dependence of $\vec{\tilde{E}}$ & $\vec{\tilde{H}}$ are on z only, i.e.,

$$\vec{\tilde{E}}(space) = \vec{\tilde{E}}(z) = \vec{e}_x \hat{E}_x(z) + \vec{e}_y \hat{E}_y(z) + \vec{e}_z \hat{E}_z(z)$$

$$\vec{\tilde{H}}(space) = \vec{\tilde{H}}(z) = \vec{e}_x \hat{H}_x(z) + \vec{e}_y \hat{H}_y(z) + \vec{e}_z \hat{H}_z(z)$$

deduce the following equations :

$$\hat{E}_z(z) = \text{constant} \quad \dots\dots (9)$$

$$\hat{H}_z(z) = \text{constant} \quad \dots\dots (10)$$

$$\frac{d \hat{E}_y(z)}{dz} = i\omega\mu \hat{H}_x(z) \quad \dots\dots (11)$$

$$\frac{d \hat{E}_x(z)}{dz} = -i\omega\mu \hat{H}_y(z) \quad \dots\dots (12)$$

$$\hat{H}_z(z) = 0 \quad \dots\dots (13)$$

$$\frac{d \hat{H}_y(z)}{dz} = -(\sigma + i\omega\varepsilon) \hat{E}_x(z) \quad \dots\dots (14)$$

$$\frac{d \hat{H}_x(z)}{dz} = (\sigma + i\omega\varepsilon) \hat{E}_y(z) \quad \dots\dots (15)$$

$$\hat{E}_z(z) = 0 \quad \dots\dots (16)$$

(10 marks)

Question three (continued)

- (c) From equations in (b) deduce the following wave equation for $\hat{E}_y(z)$ as

$$\frac{d^2 \hat{E}_y(z)}{dz^2} = i\omega\mu(\sigma + i\omega\varepsilon)\hat{E}_y(z) \dots\dots (17)$$

then by direct substitution show that

$$\hat{E}_y(z) = \hat{E}_m^+ e^{-\hat{\gamma}z} + \hat{E}_m^- e^{+\hat{\gamma}z} \dots (18) \quad \text{where } \hat{\gamma} = \sqrt{i\omega\mu(\sigma + i\omega\varepsilon)} \text{ and}$$

\hat{E}_m^+ & \hat{E}_m^- are complex constants (represents + & - z complex amplitudes)

is solution to eq(17).

(7 marks)

- (d) Substitute eq(18) into one of the equations in (b) and deduce that its wave partner $\hat{H}_x(z)$ is

$$\hat{H}_x(z) = \hat{H}_m^+ e^{-\hat{\gamma}z} + \hat{H}_m^- e^{+\hat{\gamma}z} \dots (19) \quad \text{where}$$

$$-\frac{\hat{E}_m^+}{\hat{H}_m^+} = +\frac{\hat{E}_m^-}{\hat{H}_m^-} = \sqrt{\frac{i\omega\mu}{\sigma + i\omega\varepsilon}} \xrightarrow{\text{set as}} \hat{\eta}$$

(5 marks)

Question four

(a) Given the following data for Cesium metal Cs as :

Atomic weight = 132.9 gm / gm-mole

Density = 1873 kg / m³

Conductivity $\sigma = 4.9 \times 10^7 \frac{1}{\Omega m}$ at room temperature

(i) Find the number of conduction electrons per meter cube, i.e., number density n , for metal Cs if each Cs atom has one conduction electron. **(4 marks)**

(Hint : one gm – mole pure metal contains 6.022×10^{23} atoms)

(ii) Find the value of the mean free time t_f for metal Cs at room temperature.

(Hint : $t_f = \frac{2 m_e \sigma}{n e^2}$) **(4 marks)**

(iii) If the applied constant electric field is $\vec{E} = \vec{a}_y 100 \text{ V/m}$, find the saturated drifting velocity \vec{v}_d of the average conducting electron of metal Cs at room temperature. **(4 marks)**

(Hint : $\vec{v}_d = -\frac{e t_f}{2 m_e} \vec{E}$)

(b) An uniform plane wave travelling along +z direction with the field components of

(E_x, H_y) has a complex electric field amplitude of $100 e^{i50^\circ} \frac{V}{m}$ and propagates at

$f = 5 \times 10^6 \text{ Hz}$ in a material region having the parameters of $\mu = \mu_0$, $\epsilon = 4 \epsilon_0$ &

$\frac{\sigma}{\omega \epsilon} = 1$,

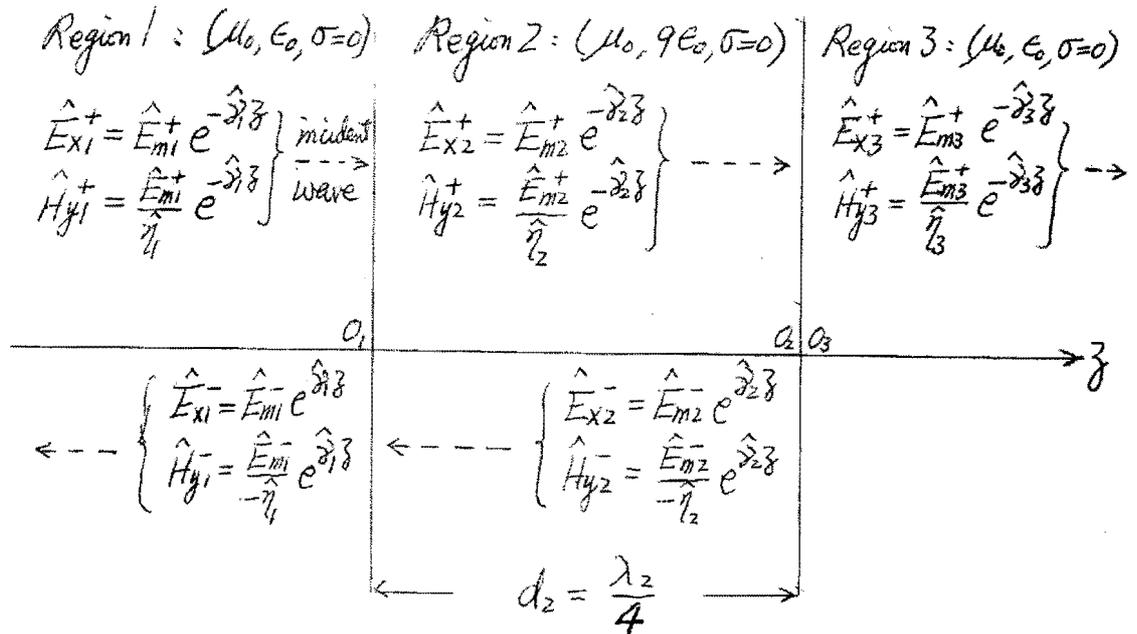
(i) Find the values of the propagation constant $\hat{\gamma} (= \alpha + i \beta)$ and the intrinsic wave impedance $\hat{\eta}$ for this wave. **(5 marks)**

(ii) express the electric and magnetic fields in both their complex and real-time forms, with the numerical values of (b)(i) inserted, **(5 marks)**

(iii) find the values of the penetration depth, wave length and phase velocity of the given wave. **(3 marks)**

Question five

An uniform plane wave $(\hat{E}_{x1}^+, \hat{H}_{y1}^+)$, operates at $f = 10^8$ Hz, is normally incident upon a lossless layer of $d_2 \left(= \frac{\lambda_2}{4} \right)$ thickness with parameters of $(\mu = \mu_0, \epsilon = 9 \epsilon_0)$ as shown below :



0_1 , 0_2 & 0_3 are the respective origins for region 1, 2 & 3 chosen at the first and second interface. (Both region 1 and region 3 are air region.)

- (a) Find the values of $\hat{\gamma}_1$, $\hat{\gamma}_2$, $\hat{\gamma}_3$, $\hat{\eta}_2$ & λ_2 , (note: $\hat{\eta}_1 = \hat{\eta}_3 = 120 \pi \Omega$ and $\alpha_1 = \alpha_2 = \alpha_3 = 0$) (4 marks)
- (b) Starting with $\hat{\Gamma}_3(z) = 0$ for the rightmost region, i.e., region 3, and using continuous \hat{Z} at the interface as well as the equations in (a), find the values of $\hat{Z}_3(0)$, $\hat{Z}_2(0)$, $\hat{\Gamma}_2(0)$, $\hat{\Gamma}_2(-d_2)$, $\hat{Z}_2(-d_2)$, $\hat{Z}_1(0)$ & $\hat{\Gamma}_1(0)$ (10 marks)
- (c) Find the value of \hat{E}_{m1}^- , \hat{E}_{m2}^+ , \hat{E}_{m2}^- & \hat{E}_{m3}^+ if given $\hat{E}_{m1}^+ = 50 e^{i40^\circ} \frac{V}{m}$. (11 marks)

Useful informations

$$e = 1.6 \times 10^{-19} \text{ C}$$

$$m_e = 9.1 \times 10^{-31} \text{ kg}$$

$$\mu_0 = 4 \pi \times 10^{-7} \frac{\text{H}}{\text{m}}$$

$$\varepsilon_0 = 8.85 \times 10^{-12} \frac{\text{F}}{\text{m}}$$

$$\alpha = \frac{\omega \sqrt{\mu \varepsilon}}{\sqrt{2}} \sqrt{\sqrt{1 + \left(\frac{\sigma}{\omega \varepsilon}\right)^2} - 1}$$

$$\beta = \frac{\omega \sqrt{\mu \varepsilon}}{\sqrt{2}} \sqrt{\sqrt{1 + \left(\frac{\sigma}{\omega \varepsilon}\right)^2} + 1}$$

$$\frac{1}{\sqrt{\mu_0 \varepsilon_0}} = 3 \times 10^8 \frac{\text{m}}{\text{s}}$$

$$\hat{\eta} = \frac{\sqrt{\frac{\mu}{\varepsilon}}}{\sqrt{1 + \left(\frac{\sigma}{\omega \varepsilon}\right)^2}} e^{i \frac{1}{2} \tan^{-1} \left(\frac{\sigma}{\omega \varepsilon}\right)}$$

$$\eta_0 = 120 \pi \Omega = 377 \Omega$$

$$\beta_0 = \omega \sqrt{\mu_0 \varepsilon_0}$$

$$\oiint_S \vec{E} \cdot d\vec{s} = \frac{1}{\varepsilon} \iiint_V \rho_v dv$$

$$\oiint_S \vec{B} \cdot d\vec{s} = 0$$

$$\oint_L \vec{E} \cdot d\vec{l} = -\frac{\partial}{\partial t} \left(\iint_S \vec{B} \cdot d\vec{s} \right)$$

$$\oint_L \vec{B} \cdot d\vec{l} = \mu \iint_S \vec{J} \cdot d\vec{s} + \mu \varepsilon \frac{\partial}{\partial t} \left(\iint_S \vec{E} \cdot d\vec{s} \right)$$

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho_v}{\varepsilon}$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \times \vec{B} = \mu \vec{J} + \mu \varepsilon \frac{\partial \vec{E}}{\partial t}$$

$$\vec{J} = \sigma \vec{E}$$

$$\oiint_S \vec{F} \cdot d\vec{s} \equiv \iiint_V (\vec{\nabla} \cdot \vec{F}) dV \quad \text{divergence theorem}$$

$$\oint_L \vec{F} \cdot d\vec{l} \equiv \iint_S (\vec{\nabla} \times \vec{F}) \cdot d\vec{s} \quad \text{Stokes' theorem}$$

$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{F}) \equiv 0$$

$$\vec{\nabla} \times (\vec{\nabla} f) \equiv 0$$

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{F}) \equiv \vec{\nabla} (\vec{\nabla} \cdot \vec{F}) - \nabla^2 \vec{F}$$

$$\begin{aligned} \vec{\nabla} f &= \vec{e}_x \frac{\partial f}{\partial x} + \vec{e}_y \frac{\partial f}{\partial y} + \vec{e}_z \frac{\partial f}{\partial z} = \vec{e}_\rho \frac{\partial f}{\partial \rho} + \vec{e}_\phi \frac{1}{\rho} \frac{\partial f}{\partial \phi} + \vec{e}_z \frac{\partial f}{\partial z} \\ &= \vec{e}_r \frac{\partial f}{\partial r} + \vec{e}_\theta \frac{1}{r} \frac{\partial f}{\partial \theta} + \vec{e}_\phi \frac{1}{r \sin(\theta)} \frac{\partial f}{\partial \phi} \end{aligned}$$

$$\begin{aligned} \vec{\nabla} \cdot \vec{F} &= \frac{\partial(F_x)}{\partial x} + \frac{\partial(F_y)}{\partial y} + \frac{\partial(F_z)}{\partial z} = \frac{1}{\rho} \frac{\partial(F_\rho \rho)}{\partial \rho} + \frac{1}{\rho} \frac{\partial(F_\phi)}{\partial \phi} + \frac{\partial(F_z)}{\partial z} \\ &= \frac{1}{r^2} \frac{\partial(F_r r^2)}{\partial r} + \frac{1}{r \sin(\theta)} \frac{\partial(F_\theta \sin(\theta))}{\partial \theta} + \frac{1}{r \sin(\theta)} \frac{\partial(F_\phi)}{\partial \phi} \end{aligned}$$

$$\begin{aligned} \vec{\nabla} \times \vec{F} &= \vec{e}_x \left(\frac{\partial(F_z)}{\partial y} - \frac{\partial(F_y)}{\partial z} \right) + \vec{e}_y \left(\frac{\partial(F_x)}{\partial z} - \frac{\partial(F_z)}{\partial x} \right) + \vec{e}_z \left(\frac{\partial(F_y)}{\partial x} - \frac{\partial(F_x)}{\partial y} \right) \\ &= \frac{\vec{e}_\rho}{\rho} \left(\frac{\partial(F_z)}{\partial \phi} - \frac{\partial(F_\phi \rho)}{\partial z} \right) + \vec{e}_\phi \left(\frac{\partial(F_\rho)}{\partial z} - \frac{\partial(F_z)}{\partial \rho} \right) + \frac{\vec{e}_z}{\rho} \left(\frac{\partial(F_\phi \rho)}{\partial \rho} - \frac{\partial(F_\rho)}{\partial \phi} \right) \\ &= \frac{\vec{e}_r}{r^2 \sin(\theta)} \left(\frac{\partial(F_\phi r \sin(\theta))}{\partial \theta} - \frac{\partial(F_\theta r)}{\partial \phi} \right) + \frac{\vec{e}_\theta}{r \sin(\theta)} \left(\frac{\partial(F_r)}{\partial \phi} - \frac{\partial(F_\phi r \sin(\theta))}{\partial r} \right) + \frac{\vec{e}_\phi}{r} \left(\frac{\partial(F_\theta r)}{\partial r} - \frac{\partial(F_r)}{\partial \theta} \right) \end{aligned}$$

where $\vec{F} = \vec{e}_x F_x + \vec{e}_y F_y + \vec{e}_z F_z = \vec{e}_\rho F_\rho + \vec{e}_\phi F_\phi + \vec{e}_z F_z = \vec{e}_r F_r + \vec{e}_\theta F_\theta + \vec{e}_\phi F_\phi$ and
 $d\vec{l} = \vec{e}_x dx + \vec{e}_y dy + \vec{e}_z dz = \vec{e}_\rho d\rho + \vec{e}_\phi \rho d\phi + \vec{e}_z dz = \vec{e}_r dr + \vec{e}_\theta r d\theta + \vec{e}_\phi r \sin(\theta) d\phi$

$$\begin{aligned} \nabla^2 f &= \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2} = \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial f}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 f}{\partial \phi^2} + \frac{\partial^2 f}{\partial z^2} \\ &= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial f}{\partial r} \right) + \frac{1}{r^2 \sin(\theta)} \frac{\partial}{\partial \theta} \left(\sin(\theta) \frac{\partial f}{\partial \theta} \right) + \frac{1}{r^2 \sin^2(\theta)} \frac{\partial^2 f}{\partial \phi^2} \end{aligned}$$

Drude's model $m_e \frac{d\vec{v}_d}{dt} = -e \vec{E} - \frac{2 m_e \vec{v}_d}{t_f}$