UNIVERSITY OF SWAZILAND

FACULTY OF SCIENCE

DEPARTMENT OF PHYSICS

MAIN EXAMINATION 2005

TITLE OF PAPER :

ELECTROMAGNETIC THEORY

COURSE NUMBER

P331

:

:

TIME ALLOWED

THREE HOURS

INSTRUCTIONS

ANSWER ANY FOUR OUT OF FIVE

QUESTIONS.

EACH QUESTION CARRIES 25 MARKS.

MARKS FOR DIFFERENT SECTIONS ARE SHOWN IN THE RIGHT-HAND MARGIN.

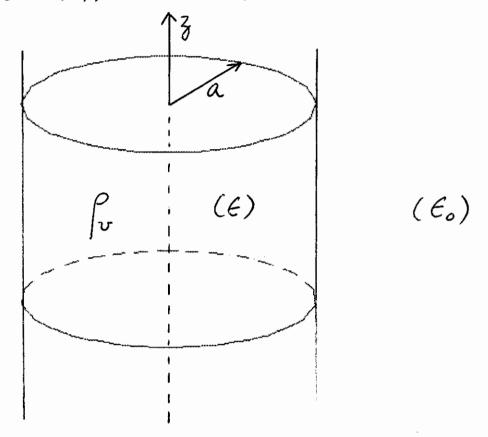
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P331 Electromagnetic Theory I

Question one

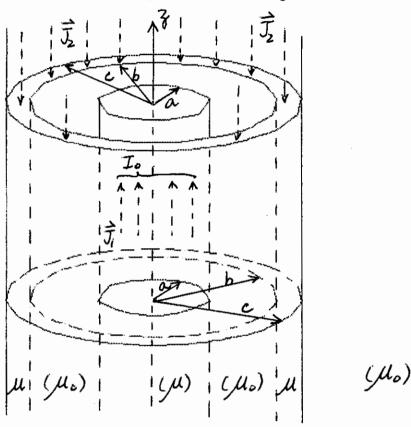
(a) For a very long wire of cross sectional radius a, if uniformly charged, i.e., the volume charge density $\rho_v = \text{constant}$ for $0 \le \rho \le a$, as shown below,



use the integral form of Coulomb's Law , i.e., $\oint_S \vec{D} \bullet d\vec{s} = \iiint_V \rho_v \, dv$, to find the electric field \vec{E} in the $0 \le \rho \le a$ (material region with permittivity ε) and $\rho \ge a$ (free space with permittivity ε_0) regions respectively, (8 marks)

Question one (continued)

(b) A very long straight coaxial cable, with an inner solid wire of radius a carries an uniform current density $\vec{J}_1 = \vec{a}_z \, \frac{I_0}{\pi \, a^2}$ and outer hollow wire with inner radius b and outer radius c carries an uniform current density $\vec{J}_2 = \vec{a}_z \, \frac{-I_0}{\pi \, (c^2 - b^2)}$ where I_0 is the total current in the inner wire, as shown in the diagram below:



(i) Use the integral form of Ampere's Law , i.e., $\oint_{l} \vec{H} \cdot d\vec{l} = \iint_{\mathcal{S}} \vec{J} \cdot d\vec{s}$, to find the magnetic field \vec{B} for $0 \le \rho \le a$ (material region with permeability μ), $a \le \rho \le b$ (free space with permeability μ_0), $b \le \rho \le c$ (material region with permeability μ) and $\rho \ge c$ (free space with permeability μ_0) regions respectively,

Question one (continued)

(ii) Find the total magnetic energy, i.e., $U_m = \frac{1}{2} \iiint_V (\vec{B} \cdot \vec{H}) dv$, stored in the given cable of length l between the wires, i.e., $a \le \rho \le b$, and then find the external self-inductance L_e for this length of coaxial cable. (5 marks)

(Note:
$$U_m = \frac{1}{2} L_e I_0^2$$
 and $\iiint_V dv = \int_{\rho = a}^b \int_{\phi = 0}^{2\pi} \int_{z=0}^l \rho \, d\rho \, d\phi \, dz$)

Question two

(a) Write down the continuity equation of electric charge in its integral and differential forms. $\vec{a} \vec{D}$

Show that without introducing the displacement current term ,i.e., $\frac{\partial D}{\partial t}$, in the equation for Ampere's law , i.e., $\vec{\nabla} \times \vec{H} = \vec{J}$ instead of $\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$, Maxwell's

equations would be in contradiction with the continuity equation . (8 marks)

(b) In a conductive region, based on Drude model, the force on a conducting electron by \vec{E} is $-e\vec{E}$ and the average momentum change caused by the ion lattice of the conductor is $-\frac{m_e\vec{v}_d}{\tau}$ and, thus, in a steady state situation one has $-e\vec{E}-\frac{m_e\vec{v}_d}{\tau}=0$,

(i) Explain briefly \vec{v}_d and τ_c . (2 marks)

- (ii) Using the current density in the conductive region $\vec{J} = \rho_v \ \vec{v}_d = -n \ e \ \vec{v}_d$, deduce the following point form of Ohm's Law $\vec{J} = \sigma \ \vec{E}$ where $\sigma = \frac{n \ e^2}{m_e} \tau_c$ for the dc case.
- (iii) Sodium has an atomic density of $2.3 \times 10^{28} \frac{atom}{m^3}$ at room temperature, and with one outer-orbit electron available for conduction, find the value of τ_c for sodium if its dc conductivity $\sigma = 2.1 \times 10^7 \frac{1}{m\Omega}$ at room temperature.

(3 marks)

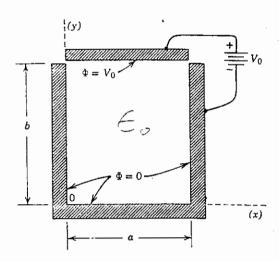
Question two (continued)

- (c) For the time-independent (or static) case , i.e., $\frac{\partial}{\partial t} \to 0$, setting $\vec{E} = -\vec{\nabla} \Phi$ and $\vec{B} = \vec{\nabla} \times \vec{A}$, where Φ and \vec{A} are the scalar and vector potentials respectively ,
 - (i) show that Maxwell's equations for Faraday's law and for the magnetic Coulomb's law are automatically satisfied. (3 marks)
 - (ii) Use Coulomb's gauge ,i.e., $\vec{\nabla} \cdot \vec{A} = 0$, to deduce the following equations :

$$\nabla^2 \vec{A} = -\mu \vec{J}$$
 and $\nabla^2 \Phi = -\frac{\rho_v}{\varepsilon}$. (6 marks)

Question three

An infinitely long, rectangular U shaped conducting channel is insulated at the corners from a conducting plate forming the fourth side with interior dimensions as shown below:



The differential equation for the electric potential $\Phi(x, y)$ in between the two conductors is

$$\frac{\partial^{2}\Phi\left(x,y\right)}{\partial x^{2}} + \frac{\partial^{2}\Phi\left(x,y\right)}{\partial y^{2}} = 0$$

(a) Setting $\Phi(x,y) = X(x)Y(y)$ and applying the technique of separation of variables, deduce the following two ordinary differential equations:

$$\begin{cases} \frac{d^2 X(x)}{d x^2} = -k^2 X(x) \cdot \dots \cdot (1) \\ \frac{d^2 Y(y)}{d y^2} = k^2 Y(y) \cdot \dots \cdot (2) \end{cases}$$

where k (> 0) is a separation constant.

(3 marks)

Question three (continued)

- (b) For a particular value of k, show that the general solution for equations (1) & (2) can be written as: $X_k(x) = (A_k \cos(kx) + B_k \sin(kx))$, $Y_k(y) = (C_k \cosh(ky) + D_k \sinh(ky))$ and $\Phi_k(x,y) = X_k(x) Y_k(y)$ where A_k , B_k , C_k and D_k are arbitrary constants. (3 marks)
- (c) Applying the given boundary conditions for Φ at x=0, y=0 and x=a, i.e., $\Phi_k(0,y)=\Phi_k(a,y)=\Phi_k(x,0)=0$, deduce that $A_k=0$, $C_k=0 \quad \text{and} \quad k=\frac{n\pi}{a} \quad , \quad n=1,2,\cdots \qquad . \tag{6 marks}$
- (d) The general solution now can be written as $\Phi(x, y) = \sum_{n=1}^{\infty} \Phi_n(x, y)$ where

$$\Phi_n(x,y) = E_n \sinh(\frac{n\pi}{a}y) \sin(\frac{n\pi}{a}x)$$
 , (k re-indexed as n and $B_n D_n \equiv E_n$)

Apply the last boundary condition at y = b, i.e., $\Phi(x,b) = V_0$, utilize the Fourier

integral
$$\int_0^a \sin(\frac{n\pi}{a}x) \sin(\frac{m\pi}{a}x) dx = \begin{cases} 0 & \text{if } n \neq m \\ \frac{a}{2} & \text{if } n = m \end{cases}$$
 to find

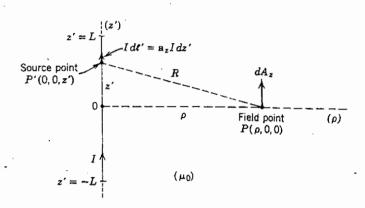
 $E_n = 1, 2, 3, \dots$. Show that the specific solution for this problem is

$$\Phi(x,y) = \frac{4V_0}{\pi} \left(\frac{\sinh(\frac{\pi y}{a})}{\sinh(\frac{\pi b}{a})} \sin(\frac{\pi x}{a}) + \frac{1}{3} \frac{\sinh(\frac{3\pi y}{a})}{\sinh(\frac{3\pi b}{a})} \sin(\frac{3\pi x}{a}) + \cdots \right)$$

(13 marks)

Question four

(a) A straight thin wire of finite length 2L in free space carries a direct current I as shown in the diagram below:



(i) For a field point $P(\rho,0,0)$ on a plane bisecting the given conducting wire,

use
$$\vec{A}(\rho,0,0) = \int_{l} \frac{\mu_0 I d\vec{l}'}{4 \pi R}$$
 where $d\vec{l}' = \vec{a}_z dz'$,

$$z' = -L$$
 to L and $R = \sqrt{\rho^2 + (z')^2}$, to deduce that

$$\vec{A} = \vec{a}_z \frac{\mu_0 I}{4 \pi} \ln \left(\frac{\sqrt{L^2 + \rho^2} + L}{\sqrt{L^2 + \rho^2} - L} \right)$$
 (6 marks)

(Hint:
$$\int \frac{dx}{\sqrt{a^2 + x^2}} = \ln(x + \sqrt{a^2 + x^2})$$
)

(ii) Find the magnetic field \vec{B} from the vector potential \vec{A} obtained in (a)(i).

Deduce that $\vec{B} = \vec{a}_{\phi} \frac{\mu_0 I}{2 \pi \rho} \frac{L}{\sqrt{L^2 + \rho^2}}$ and thus show that if $L >> \rho$ then

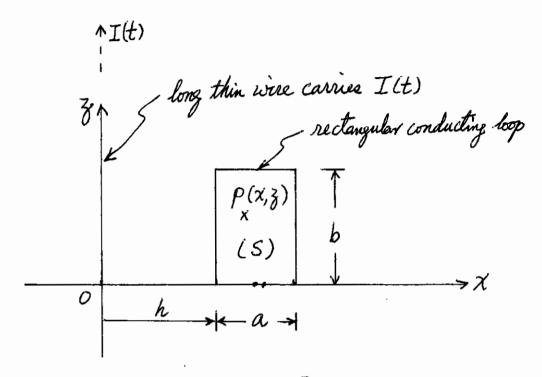
$$\vec{B} \approx \vec{a}_{\phi} \, \frac{\mu_0 \, I}{2 \, \pi \, \rho}$$
 which is the \vec{B} field generated by an infinitely long current

carrying wire.

(7 marks)

Question four (continued)

(b) A rectangular conducting loop is situated near a very long thin straight wire carrying a slow time-variant current I(t) as show in the diagram below:



- (i) Use the quasi-static expression of \vec{B} in (a)(ii) for an infinitely long current carrying wire with I replaced by I(t), \vec{a}_{ϕ} replaced by \vec{a}_{y} and ρ replaced by x, i.e., $\vec{B} = \vec{a}_{y} \frac{\mu_{0} I(t)}{2 \pi x}$, to find the total \vec{B} flux lines out of the given rectangular conducting loop, i.e., $\iint_{S} \vec{B} \cdot d\vec{s}$ where $d\vec{s} = -\vec{a}_{y} dx dz$, with x = h to h + a and z = 0 to b. (8 marks)
- (ii) Use Faraday's Law, find the induced e.m.f. on the given rectangular conducting loop if $I(t) = I_0 \cos(\omega t)$. (4 marks)

Question five

(a) (i) If \vec{E} and \vec{H} are replaced by $\hat{\vec{E}}(\vec{r}) e^{i\omega t}$ and $\hat{\vec{H}}(\vec{r}) e^{i\omega t}$ respectively, deduce the following time harmonic Maxwell's equations for a material region represented by parameters of μ , ε , σ and also having $\rho_{\nu} = 0$ as:

$$\vec{\nabla} \cdot \vec{\hat{E}}(\vec{r}) = 0 \quad , \quad \vec{\nabla} \cdot \vec{\hat{H}}(\vec{r}) = 0 \quad , \quad \vec{\nabla} \times \vec{\hat{E}}(\vec{r}) = -i\omega \ \mu \ \vec{\hat{H}}(\vec{r}) \quad \text{and}$$

$$\vec{\nabla} \times \vec{\hat{H}}(\vec{r}) = \left(\sigma + i\omega \ \varepsilon\right) \vec{\hat{E}}(\vec{r}) \qquad (4 \text{ marks})$$

(ii) If further assuming that $\ \vec{\hat{E}}$ and $\ \vec{\hat{H}}$ are functions of z only , i.e.,

$$\vec{\hat{E}}(\vec{r}) = \vec{e}_x \ \hat{E}_x(z) + \vec{e}_y \ \hat{E}_y(z) + \vec{e}_z \ \hat{E}_z(z) \qquad \text{and} \qquad$$

$$\vec{\hat{H}}(\vec{r}) = \vec{e}_x \, \hat{H}_x(z) + \vec{e}_y \, \hat{H}_y(z) + \vec{e}_z \, \hat{H}_z(z) \ , \label{eq:hamiltonian}$$

deduce the following equations:

$$\hat{E}_z(z) = 0 = \hat{H}_z(z) \quad , \quad \frac{d\,\hat{E}_x(z)}{d\,z} = -\,i\,\omega\,\,\mu\,\hat{H}_y(z) \quad , \label{eq:energy}$$

$$\frac{d\hat{E}_{y}(z)}{dz} = i\omega \,\mu \,\hat{H}_{x}(z) \quad , \quad \frac{d\hat{H}_{x}(z)}{dz} = (\sigma + i\omega \,\varepsilon) \hat{E}_{y}(z) \quad ,$$

$$\frac{d\hat{H}_{y}(z)}{dz} = -(\sigma + i\omega \varepsilon)\hat{E}_{x}(z) \quad \text{and} \quad$$

$$\frac{d^2 \hat{E}_x(z)}{dz^2} + (\omega^2 \mu \varepsilon - i \omega \mu \sigma) \hat{E}_x(z) = 0$$
 (8 marks)

Question five (continued)

- (b) An uniform plane wave travelling along +z direction with the field components $E_x(z) \quad \text{and} \quad H_y(z) \quad \text{has an complex electric field amplitude of} \quad 100 \, e^{i\,30^\circ} \, \frac{V}{m} \quad \text{and}$ propagates at $f=10^6 \, Hz$ in a material region having the parameters $\mu=\mu_0$, , $\varepsilon=9\,\varepsilon_0 \quad \text{and} \quad \frac{\sigma}{\omega\,\varepsilon}=0.5 \quad .$
 - (i) Find the values of the propagation constant $\hat{\gamma}$ (= α + $i\beta$) and the intrinsic wave impedance $\hat{\eta}$ for this wave. (4 marks)
 - (ii) Express the electric and magnetic fields in both their complex and real-time forms, with the numerical values of (b)(i) inserted. (4 marks)
 - (iii) Find the values of the penetration depth, wave length and phase velocity of the given wave. (5 marks)

USEFUL INFORMATION

$$e = 1.6 \times 10^{-19}$$
 C , $m_e = 9.1 \times 10^{-31}$ kg

$$\frac{1}{\sqrt{\mu_0 \, \varepsilon_0}} = 3 \times 10^8 \, \frac{m}{s} \qquad , \qquad \beta_0 \equiv \omega \, \sqrt{\mu_0 \, \varepsilon_0} \qquad , \quad \mu_0 = 4 \, \pi \times 10^{-7} \, \frac{H}{m}$$

$$, \varepsilon_0 = 8.85 \times 10^{-12} \frac{F}{m} \qquad , \qquad \hat{\eta} = \frac{\sqrt{\frac{\mu}{\varepsilon}}}{\sqrt{1 + \left(\frac{\sigma}{\omega \, \varepsilon}\right)^2}} e^{i\frac{1}{2}\tan^{-1}\left(\frac{\sigma}{\omega \, \varepsilon}\right)}$$

$$\eta_0 = \sqrt{\frac{\mu_0}{\varepsilon_0}} = 120 \pi \quad \Omega = 377 \quad \Omega$$

$$\alpha = \frac{\omega \sqrt{\mu \, \varepsilon}}{\sqrt{2}} \sqrt{\sqrt{1 + \left(\frac{\sigma}{\omega \, \varepsilon}\right)^2} - 1} \qquad , \qquad \beta = \frac{\omega \sqrt{\mu \, \varepsilon}}{\sqrt{2}} \sqrt{\sqrt{1 + \left(\frac{\sigma}{\omega \, \varepsilon}\right)^2} + 1}$$

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\varepsilon} \quad , \quad \vec{\nabla} \cdot \vec{B} = 0 \quad , \quad \vec{\nabla} \times \vec{E} = - \, \frac{\partial \, \vec{B}}{\partial \, t} \quad , \quad \vec{\nabla} \times \vec{B} = \mu \, \vec{J} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{\nabla} \times \vec{B} = \mu \, \vec{J} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{B} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{B} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{B} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{B} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{B} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{B} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{B} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{B} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{B} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial \, \vec{E}}{\partial \, t} \quad , \quad \vec{D} \times \vec{D} = \mu \, \vec{D} + \mu \, \varepsilon \, \frac{\partial$$

$$\vec{\nabla} \times \vec{B} = \mu \; \sigma \; \vec{E} + \mu \; \varepsilon \; \frac{\partial \vec{E}}{\partial t} \; , \label{eq:delta_energy}$$

$$\vec{\nabla}\times\left(\vec{\nabla}\times\vec{F}\right)\equiv\vec{\nabla}\left(\vec{\nabla}\cdot\vec{F}\right)-\nabla^{2}\vec{F}\qquad,\qquad \oiint_{S}\vec{F}\cdot d\vec{s}=\iiint_{V}\vec{\nabla}\cdot\vec{F}\ dv\qquad,$$

$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{F}) = 0 \quad , \quad \vec{\nabla} \times (\vec{\nabla} f) = 0$$

$$\vec{\nabla} f = \vec{e}_r \frac{\partial f}{\partial r} + \vec{e}_\theta \frac{1}{r} \frac{\partial f}{\partial \theta} + \vec{e}_\phi \frac{1}{r \sin \theta} \frac{\partial f}{\partial \phi} \qquad ,$$

$$\vec{\nabla} \bullet \vec{F} = \frac{\partial F_x}{\partial x} + \frac{\partial F_y}{\partial y} + \frac{\partial F_z}{\partial z}$$

USEFUL INFORMATION (continued)

$$\vec{\nabla} \cdot \vec{F} = \frac{1}{r^2} \frac{\partial \left(r^2 F_r\right)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial \left(\sin \theta F_\theta\right)}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial F_\phi}{\partial \phi} \quad ,$$

$$\vec{\nabla} \times \vec{F} = \vec{e}_x \left[\frac{\partial F_z}{\partial y} - \frac{\partial F_y}{\partial z} \right] + \vec{e}_y \left[\frac{\partial F_x}{\partial z} - \frac{\partial F_z}{\partial x} \right] + \vec{e}_z \left[\frac{\partial F_y}{\partial x} - \frac{\partial F_x}{\partial y} \right] \quad ,$$

$$\vec{\nabla} \times \vec{F} = \frac{\vec{e}_r}{r \sin \theta} \left[\frac{\partial}{\partial \theta} \left(F_{\phi} \sin \theta \right) - \frac{\partial F_{\theta}}{\partial \phi} \right] + \frac{\vec{e}_{\theta}}{r} \left[\frac{1}{\sin \theta} \frac{\partial F_r}{\partial \phi} - \frac{\partial \left(r F_{\phi} \right)}{\partial r} \right] + \frac{\vec{e}_{\phi}}{r} \left[\frac{\partial \left(r F_{\theta} \right)}{\partial r} - \frac{\partial F_r}{\partial \theta} \right] ,$$

$$\vec{\nabla} \times \vec{F} = \vec{a}_{\rho} \left[\frac{1}{\rho} \frac{\partial F_{z}}{\partial \phi} - \frac{\partial F_{\phi}}{\partial z} \right] + \vec{a}_{\phi} \left[\frac{\partial F_{\rho}}{\partial z} - \frac{\partial F_{z}}{\partial \rho} \right] + \vec{a}_{z} \left[\frac{1}{\rho} \frac{\partial (\rho F_{\phi})}{\partial \rho} - \frac{1}{\rho} \frac{\partial F_{\rho}}{\partial \phi} \right]$$

$$\nabla^2 f = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2}$$